

The recurrence time of Dansgaard-Oeschger events and limits on the possible periodic component

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By comparing the high-resolution isotopic records from the GRIP and NGRIP icecores, we approximately separate the climate signal from local noise to obtain an objective criterion for defining Dansgaard-Oeschger events. Our analysis identifies several additional short lasting events, increasing the total number of DO events to 27 in the period 12-90 kyr BP. The quasi-regular occurrence of the DO events could indicate a stochastic or coherent resonance mechanism governing their origin. From the distribution of waiting times we obtain a statistical upper bound on the strength of a possible periodic forcing. This finding indicates that the climate shifts are purely noise driven with no underlying periodicity.

Introduction

Abrupt temperature shifts between cold (stadial) states and warm (interstadial) states (Dansgaard-Oeschger events) are observed in the Greenland icecore isotope records (Dansgaard *et al.*, 1993; Grootes *et al.*, 1993). The transitions into these interstadials were abrupt and with temperature changes estimated from the paleo-record of the order of 10-15 °C occurring within decades. These two distinct quasi-stable climate states are most likely linked to different modes of the Atlantic thermohaline circulation (THC) (Broecker *et al.*, 1985; Stocker *et al.*, 1992). The circulation in the warm state was similar to the present interglacial state, while in the cold state the sinking took place at lower latitudes and to shallower depths (Alley & Clark, 1999; Ganopolski & Rahmstorf, 2001). A complete cessation of deep water formation is indicated by some models (Schmittner *et al.*, 2002).

Identification of the mechanism causing these climate shifts is of primary importance for understanding the stability and mode of operation of this component of the climate system. Comparing the high-resolution isotope records from GRIP and NGRIP (North GRIP members, 2004), we can separate the climate signal from the local variability and the glaciological noise in the icecores.

We observe that for timescales longer than approximately 30 years the two icecores are strongly correlated. Using the 30-year averaged record we thus obtain an objective measure to define the DO events. The initiations and terminations of the DO events are defined from consecutive up-crossings and down-crossings through two levels in the anomaly record. The definition of the DO events by this procedure is quite robust with respect to the choice of anomaly levels. Comparing to the original visual numbering of the DO events (Dansgaard *et al.*, 1993) we find several additional isotopic fluctuations that qualify as DO events. For example, DO event 2 consists of two closely spaced DO events, which based on the stratigraphy devised by Walker *et al.* (1999) and Björk *et al.* (1998), becomes Greenland Interstadial 'GI2a' (youngest) and 'GI2c' (oldest) separated by the stadial state 'GI2b'. Using these conventions the succeeding period is the Greenland Stadial 'GS2'. For further explanation see the figure caption (figure 2).

The isotope record from the GISP2 icecore, dated using stratigraphic methods (Meese *et al.*, 1997), shows a significant peak in the spectrum at 1470 years (Yiou *et al.*, 1997), indicating a possible periodic forcing of the climate system. Alley *et al.* (2001) proposed this to be a stochastic resonance while Timmermann *et al.* (2003) proposed a coherence resonance phenomenon. The climate system itself is dominated by strongly fluctuating, irregular, fast timescale noise, so it is highly unlikely that a strictly periodic signal would be internally generated. Only simplified models with few degrees of freedom exhibit strict cyclic behavior (Paillard & Labeyrie, 1994). In a recent study Roe and Steig (2004) notes that, 1470 yrs periodicity aside, the Antarctic Byrd record is comparable with a simple autoregressive process and the same is true for the Greenland records if an additional simple threshold rule is imposed.

A periodic component in the signal would likely be the result of a non-linear response to a weak external periodic forcing, though the origin of such a forcing has not yet been identified. By observing the waiting times between consecutive DO events it was noted that the record could be interpreted as having a preferred waiting time of 1470 years and multiples of this period, corresponding to the system skipping a few transitions but still switching in phase with the external periodicity (Alley *et al.*, 2001; Schulz, 2002a; Rahmstorf, 2003). Observing the waiting times rather than the power spectrum is advantageous in a noisy signal where the cyclicity is far from being sinusoidal. In this case a large portion of the power will be in overtones, which possibly brings the power in the peak below detection into the noise level. This is the case for the icecore signal where the transitions into the interstadials are rapid and the transitions into the stadials are more gradual.

By comparing the icecore record with stochastic resonance models we are able to estimate an upper bound on the strength of an external periodic forcing potential in comparison to the internal barrier between the two climate states. We find that with maximum likelihood the data are not a result of a process with a periodic component. Since there still remains discrepancy between the GRIP and the GISP2 datings we have performed our analysis on both timescales. The GISP2 timescale agrees well with the U/Th dating of the Chinese Hulu Cave stalagmite record (Wang *et al.*, 2001) while the dating of the French Villars Cave stalagmite record (Genty *et al.*, 2003) is somewhere in between. While we cannot reject the case of no periodic component we can

determine an upper limit of a periodic forcing such that a strength above this value can be rejected with 98% confidence. This upper limit varies slightly between the GRIP and the GISP2 timescales.

1 Defining the Dansgaard-Oeschger events

The power spectra of the GRIP and GISP2 records (figure 1) are estimated from the irregularly spaced data using the 'Redfit 3.5' routine (Schulz & Mudelsee, 2002). The appearance of the 1470-year peak in the spectrum from the GISP2 data (figure 1) is caused by the regular spacing of three consecutive DO events (numbered 5,6,7) (Schulz, 2002a). This regular spacing is not nearly as pronounced in the GRIP icecore using the ss09-sea timescale (Johnsen *et al.*, 1997), explaining why the spectral peak is not as significant in the GRIP spectrum. The discrepancy will hopefully be resolved in the near future where a more precise dating of the NGRIP icecore based on annual layers counting far back in the glacial period is expected. Levels of significance are not included since these are strongly depending on the null-hypothesis noise spectrum assumed. The regular Ornstein-Uhlenbeck process usually assumed (Schulz & Mudelsee, 2002) is not relevant, since the isotope record is far from this process. Here we shall not argue for either the GISP2 or GRIP dating, although we note that one cannot favor one in comparison to the other solely based on the strength of the 1470-year spectral peak as was done by Rahmstorf (2003).

Besides the uncertainty in dating, some discrepancy exists in defining the warming events. The 'canonical' DO events (Dansgaard *et al.*, 1993) (figure 2) are based on visual inspection. Alley *et al.* (2001) used a filtering procedure and threshold levels were adopted which produced 43 warming events. Schulz (2002a) found that DO events 9, 15, 16 fell below the chosen threshold level, while the event around 65 kyr BP between events 18 and 19 was included. Rahmstorf (2003) only analyzed events after 50 kyr BP. Here event 9 was omitted while it was argued for an event ('A') prior to the Younger Dryas (YD) and the termination of YD was included. We denote his event 'A' as 'GI1c' in accordance with Walker *et al.* (1999). In the following we will argue for a procedure of defining the warming events. The problem is two-fold: A spectral filtering must be decided and threshold values and a procedure for threshold crossings must be chosen.

1.1 The filtering

There is a general consensus that a multi-millennial high-pass filter must be applied in order to eliminate the variations due to the orbital forcing. Whether this is a spectral filter or subtraction of a long-term running mean is not significant. It is with respect to this long term mean that the anomaly is defined. The important problem is to decide for a low-pass in the other end of the spectrum. Most analyses have been performed on the approximately 200-year averaged isotope data available from the web archives. However, the spectral power in the timescale range shorter than 200 years is substantial. The effect of smoothing is that the positive extreme of a short timescale warming event will be reduced and possibly brought below a chosen threshold level.

Thus analyzing a smooth signal introduces a bias towards omitting very short DO events.

The shortest timescale fluctuations in the isotope records are dominated by glaciological noise (Ditlevsen *et al.*, 2002) and local fluctuations. The low-pass should thus be determined such that the noise is reduced to an insignificant level in comparison to the true climatic fluctuations. By comparing the high-resolution GRIP and NGRIP records we can estimate the noise level. Figure 3 shows the correlation between the GRIP and the NGRIP records in the 27 - 37 kyr BP period, where the NGRIP signal has been dated by matching the DO events 3-7 to the GRIP-ss09-sea timescale. The correlation coefficient is plotted as a function of the low-pass filter on both signals. The correlation between the raw high-resolution signals is approximately 0.5, while for the 30-year low-passes the correlation has increased to more than 0.8. The two core drilling sites are located more than 300 km apart so we conclude that the 30-year low-pass signal is representative for climate fluctuations and thus appropriate for the analysis.

1.2 The threshold crossing

The two distinct quasi-stable climate states are most clearly seen in the bimodal distribution of the dust-calcium signal (Ditlevsen, 1999), but is also apparent in the $\delta^{18}\text{O}$ signal where the warm interstadial states are 'sawtooth shaped', characterized by gradual decreases in the signal prior to jumping into the stadial state (Alley, 1998; Schulz, 2002b). Accordingly, the signal can be split into these two states and two transition states, cold to warm and warm to cold (Ditlevsen *et al.*, 2002). The initiation of the warm state is defined as the first up-crossing of the high threshold following an up-crossing of the low threshold. Similarly the initiation of the cold state is defined from the first down-crossing of the lower level following a down-crossing of the upper level (figure 4). The terminations are defined as the last up (down) crossings of the lower (higher) level prior to an up (down) crossing of the higher (lower) level. The periods between terminations and initiations are the transition periods.

The reason for this definition is that if the signal fluctuates around the higher (lower) level in one of the climatic states, this does not lead to a jump unless this follows a period where the signal was below (above) the lower (higher) level. This definition is advantageous in comparison to the definition used by Schulz (2002a). Here the signal is defined to be within the warm state whenever the value of the anomaly signal is above the higher threshold. This has the consequence that a warm period can more easily be split into more periods as perhaps happens for DO event 17 (Schulz, 2002a). The definitions applied here clearly identify whether DO event 17 should be regarded as one or two events. In our analysis it remains one DO event (GI17).

In the following the lower threshold is chosen as -1.0 permil anomaly and the higher threshold as +1.5 permil anomaly from the 10-kyr high-pass isotope signal. The asymmetry is because the climate system persists in the cold state longer than in the warm state. The result is shown in figure 2. The vertical full lines indicate the first up-crossings into the warm states. The result is quite robust in the sense that the jumping times are rather insensitive to the threshold values chosen.

2 The waiting time distribution

The transitions from the cold to the warm state are much more abrupt than the opposite transitions, so the waiting times between consecutive up-crossings are less sensitive to the choices of filter frequencies and thresholds and are thus used for the further analysis. The cumulated distribution of waiting times is plotted in figure 5, left panel. From a visual inspection the observed distribution seems to be sampled from a Poisson process, which has an exponential distribution. The mean waiting time t_m is 2.8 kyr, which is defined by the cumulated exponential distribution $P(t) = 1 - \exp(-t/t_m)$ shown as the straight line. The discrepancy with the earlier findings (Alley *et al.*, 2001; Schulz, 2002a; Rahmstorf, 2003) lies mainly in the inclusion of the shortest warming events. Whether or not the doubtful event 'GS18b' is included does not change the waiting time distribution significantly and does not alter the following statistical analysis.

3 An upper bound on the periodic component

Even though the distribution shown in figure 5, left panel, disfavors the previously claimed periodic component, the record is relatively short and, as Alley *et al.* (2001) noted, could be a realization of different possible processes. The waiting time distribution for a stochastic resonance process has a step-like structure with the first big step around the period and with exponentially smaller steps for multiples of the period (figure 5, right panel). A way of quantifying the degree of periodicity is then to calculate the root mean square difference between the distribution and the best-fit exponential distribution. Assuming the existence of a periodic component with period of 1470 years, we can then estimate the strength of the periodic forcing in comparison to the barrier for purely noise induced transition. Within the framework of the stochastic resonance model this amounts to investigating the strength of the resonance.

We have generated a series of data from a stochastic resonance model with the same mean waiting time as observed in the isotope signal but with different strength of the periodic component. For each realization of length similar to the length of the isotope signal we calculate the root mean square difference from the exponential distribution. The stochastic resonance model is described by the non-autonomous Langevin equation;

$$\frac{dT}{dt} = F(T) + A \cos(2\pi t/\tau) + \sigma\eta(t). \quad (1)$$

The first term represents the drift with two stable states separated by a potential barrier. The drift is derived from a potential $F = -dU/dT$, with $U(T) = T^4 - a_3T^3 - a_2T^2 + a_1T$, where a_1, a_2, a_3 are constants (Cessi, 1994; Ditlevsen, 1999). The second term in (1) is the periodic component with period τ and amplitude A . The third term is a white noise forcing with intensity σ .

For $|A| \geq A_c$, $A_c = -a_1 - a_2a_3/2 - a_3^3/8 + 8(a_2/6 - a_3^2/16)^{3/2}$ the system will jump between the two stable states through a hysteresis loop periodically,

while for $|A| < A_c$ the bifurcation points are not reached and the jumping must be noise induced.

A schematic of the potential and the periodic forcing is shown in figure 6. Disregarding non-exponential prefactors, the Kramer waiting time for penetrating the barrier going from the well containing the stable state a (c) to the well containing the stable state c (a) is estimated $T_{ac} \sim \exp(H_{min}/\sigma^2)$ ($T_{ca} \sim \exp(H_{max}/\sigma^2)$) where H_{min} (H_{max}) is the height of the barrier (Gardiner, 1985). The criterion for the noise intensity to obtain stochastic resonance is

$$T_{ac} \ll \tau \ll T_{ca}. \quad (2)$$

The climate state will then with high probability jump from state a to state c ; within time $\tau/2$ the potential has changed due to the periodic component and the state will with high probability jump from c to a , and so on.

By varying the strength of the periodic component, still tuning the noise intensity to the stochastic resonance, we can obtain an estimate of the amplitude of the periodic forcing component in comparison to the barrier in the system. Figure 7 shows a set of realizations of equation (1) for different strengths of the periodic forcing. Each realization is represented in a set of three figures; long-top -, left - and right panels. The strength of the stochastic resonance in the system is determined by the separation in timescales expressed in equation (2). When the noise is tuned to the resonance the criterion is $\Delta \equiv T_{ca}/T_{ac} \sim \exp[(H_{max} - H_{min})/\sigma^2] \gg 1$. The four panels correspond to $\Delta = 1.0, 1.8, 2.5, 3.4$ respectively.

Going from top to bottom we see that the periodic component emerges. In the right panels the spectral peak at $f = 2\pi/\tau$ emerges and exceeds the 99% confidence level (Crowley *et al.*, 1986) for $\Delta > 2.5$. The cumulated waiting-time distributions are shown in the left panels. The step-like structure from the periodicity emerges as Δ increases. The top set of panels ($\Delta = 1.0$) shows the pure Poisson process. A quantitative measure of the deviation from the Poisson process is the root mean square (rms) distance of the waiting time distribution from an exponential. This is similar to the measure used for the Kolmogorov-Smirnov test which can not be applied when we obtain the best fit distribution from the data. In the following this rms is denoted the 'error'. For each of the models shown in figure 7 the 'error' is calculated from a sufficiently long simulation (figure 8). The accuracy is within the size of the plotting symbols.

In order to evaluate the data series, we have simulated a large ensemble of time series generated by eq. (1) of same length as the isotope record for a given noise intensity. For each realization we calculate the root mean square distance to the exponential distribution. By this procedure the distribution of 'errors' is calculated. In this way the maximum likelihood model for generating the observed record is obtained. Figure 9 shows the error distributions for the models generating the series shown in figure 7. The arrows represent the error obtained from the GRIP and GISP2 isotope records. The maximum likelihood model has $\Delta = 1$ for the GRIP dating, while $\Delta = 1.6$ for the GISP2 dating. For the model with $\Delta = 2.2$, 2% of the realizations have an error less than the one measured for the isotope signal. We thus reject the hypothesis that the data are generated by a process with $\Delta \geq 2.2$ with 98% confidence for the GRIP dating. Correspondingly, we reject a process with $\Delta \geq 3.0$ with 98%

confidence for the GISP2 dating. The NGRIP record is presently dated by fitting to the GRIP timescale and is thus not independent.

4 Summary

The isotope record shows the jumping between the warm DO interstadial state and the cold stadial state with waiting times in the millennial timescale range. We have used an objective procedure based on correlating the GRIP and NGRIP records to decide the resolution of the climate record. A high and a low threshold for the anomalies are applied. The transitions are defined from consecutive level crossings. By this procedure we obtain the 'canonical' DO events and an additional set of short events. This happens either by splitting some of the 'canonical' events (1, 2, 15) or by defining events previously in the cold periods (18, 22). Assuming the record to be generated by a dynamics described by equation (1) the strength of a periodic component in the forcing in comparison to the strength of the barrier is expressed through the non-dimensional parameter Δ . It will be relevant to extend the present analysis to include the oldest DO events from event 22 to the newly discovered event 25 (North GRIP members, 2004) when a reliable dating is obtained for the NGRIP record. The observed record is highly probable as a realization of a purely noise driven process without a periodic component. This finding suggests that a mechanism for the sudden shifts between the two distinct climate states involving changes in the North Atlantic deep water formation. The shifts could be results of the erratic fluctuations in fresh water formation and heat transfer to the ocean surface. These fluctuations were strongest in the last glacial maximum where there is a tendency for more frequent shifts (Rahmstorf, 2003; Ditlevsen *et al.*, 1996). The occurrence of event 25 shortly after the termination of the Eem period, before substantial icevolumes have build up, suggests that the same climate modes exist in the interglacial periods where the 'warm' mode persists (Ganopolski & Rahmstorf, 2001). The reason why we do not experience DO events (except perhaps the 8.2 kyr event) in the Holocene climate could be the low intensity of fluctuations rather than the stability of the 'warm' mode versus the 'cold' mode as suggested by Ganopolski and Rahmstorf (2001).

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FIGURE CAPTIONS

- Fig. 1 The power spectra of the GRIP and GISP2 icecores. The difference in the spectra is due to the discrepancy in dating. The 1470 years peak indicated by the arrows is only pronounced in the GISP2 record.
- Fig. 2 The 10 kyr high-pass of the 30 years averaged GRIP $\delta^{18}\text{O}$ record. The numbering is the 'canonical' dating from Dansgaard et al. (1993). The notation Greenland Interstadial (GI) and Greenland Stadial (GS) is adapted from Walker et al. (1999). The numbering followed by a letter is the DO events obtained using the procedure described in the text. Again following Walker et al. (1999), when a previously unsplit interstadial is split into a sequence of shorter interstadial/stadial periods it keeps its assigned 'GI' preamble. The periods are then assigned letters 'a,b,c,...' after their numbering, where 'a' is the youngest. This means that fx. 'GI15b' is a stadial state despite the 'GI' preamble. With similar convention the stadial states (GS18, GS22) are split. One exception from our convention is event GI1, which is split according to Björk et al. (1998). Only events GI1c and GI1e qualify as DO events. The anomaly levels used for the up- and down-crossings are +1.5 permil and -1.0 permil respectively. The short event 'GS18b' at 62.4 kyr BP appears with this choice of anomaly levels even though it is doubtful whether it qualifies as a DO event. This cannot be determined solely based on the icecore records. Ideally it should be found in deep sea cores, or other records indicating if it can be related to a change in the THC.
- Fig. 3 The correlation between the GRIP and the NGRIP records as a function of the low-pass. The correlation falls off for fluctuations on timescales shorter than about 30 years. For the short timescales the glaciological noise and local fluctuations dominate, while the 30 years low-pass is taken to define a robust climate signal.
- Fig. 4 The separation between stadial – and interstadial states is done on the 30 years low-pass of the isotope signal. The periods are defined by the separation points being the first down-crossing time of a lower level (lower dashed line) after the signal has crossed up through an upper level (upper dashed line). The upper level is the 1.5 permil anomaly while the lower level is 1.0 permil negative anomaly. This defines the points marked with *a*'s and *d*'s (diamonds). The *b* and *c* points (triangles) are obtained in the same way by moving backward in time (from left to right in the plot) (Ditlevsen et al., 2002).
- Fig. 5 Left panel: The cumulated waiting time distribution between consecutive up-crossings into warm events. The probability for waiting more than τ is plotted as a function of τ . The straight line is an exponential distribution with mean waiting time of 2.8 kyr. Right panel: The cumulated waiting time distribution for a long realization of a stochastic resonance. The step-like structure deviation from a straight line exponential distribution reflects the periodicity in the signal.
- Fig. 6 The potential in the two extremal positions for $t = \pm\tau/4$

- Fig. 7 The long panels show realizations of equation (1) with $\Delta = 1.0, 1.6, 2.5, 3.4$. Δ is the dimensionless ratio of expected waiting times for jumping from the shallow – to the deep – and jumping from the deep – to the shallow well, see text for explanation. The corresponding power spectra are shown in the following right panels. The 99% significance levels is defined by the a posteriori criterion (*Crowley et al., 1986*). Using this criterion, the spectral peak at τ^{-1} is significant for $\Delta \geq 2.5$. The cumulated waiting time distributions are shown in the left panels, going from the purely noise driven Poisson process towards an almost periodically shifting signal.
- Fig. 8 The root-mean-square (rms) errors expressing the deviation from an exponential distribution for the simulations shown in figure 7. The rms errors are calculated from simulations much longer than the ones shown in figure 7
- Fig. 9 The distributions of calculated rms errors from realizations of length corresponding to the length of the observed record. The arrow shows the value obtained for the isotope signal. This value is most probably generated by a process with $\Delta = 1.0$ which is the purely noise driven process.

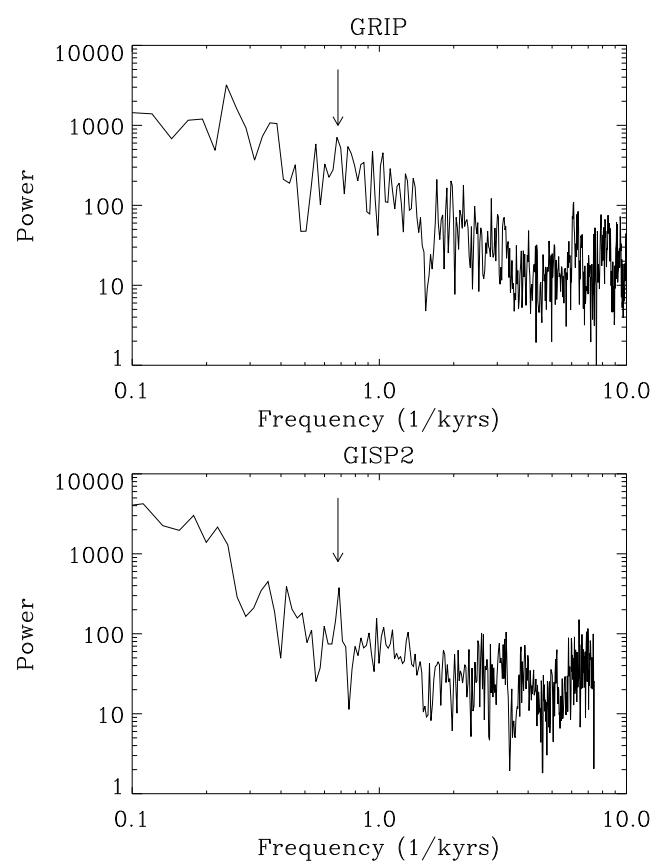


Figure 1:

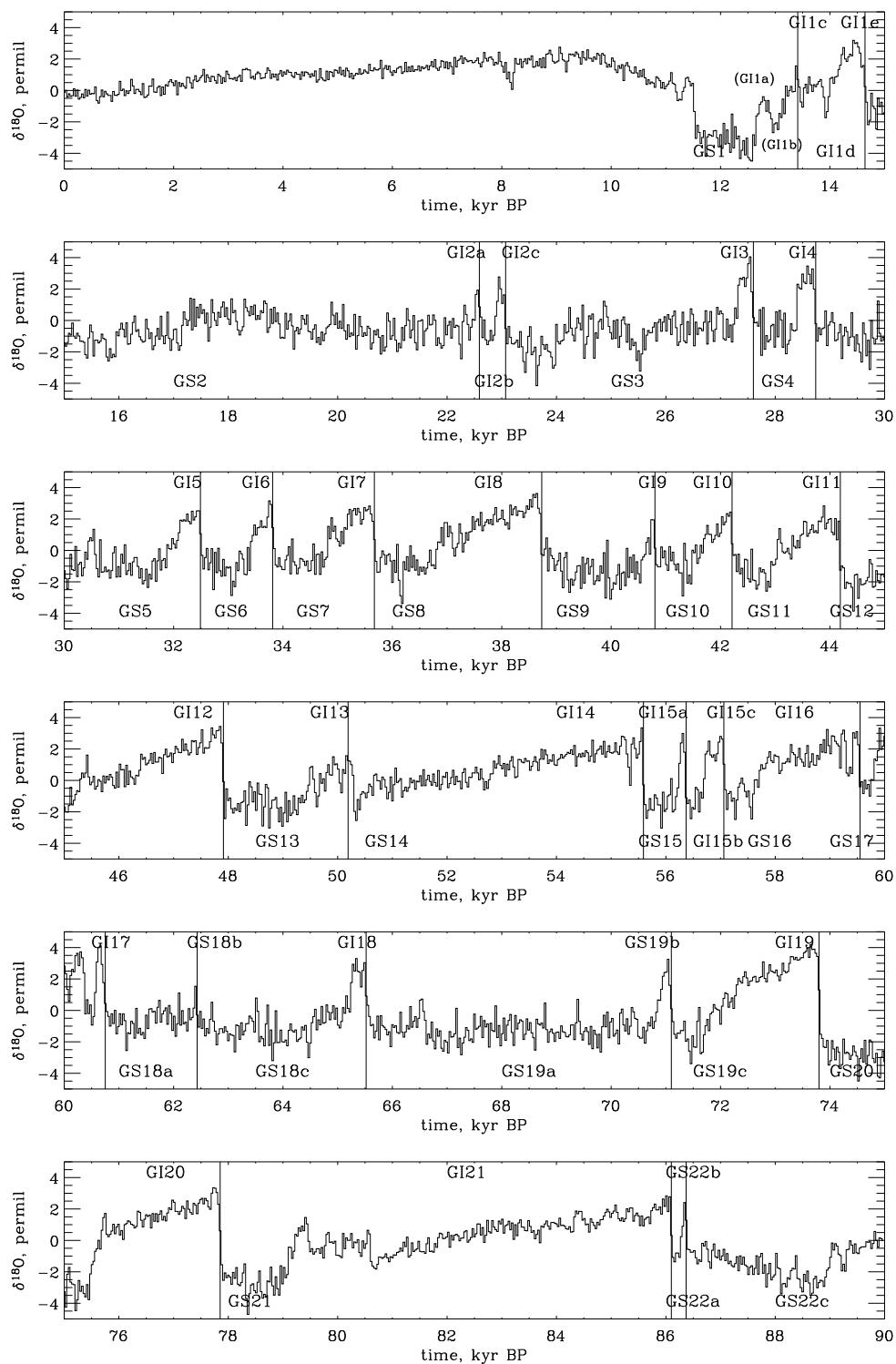


Figure 2:

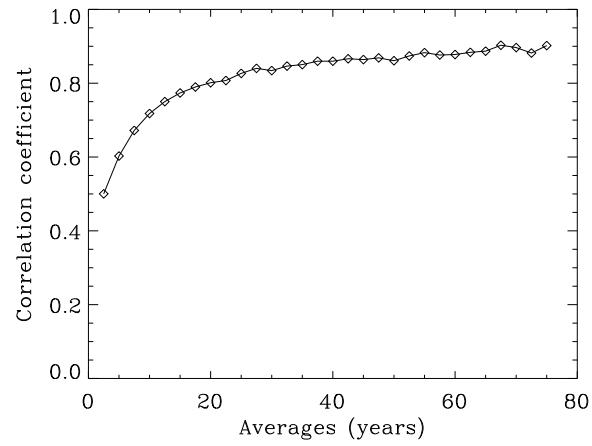


Figure 3:

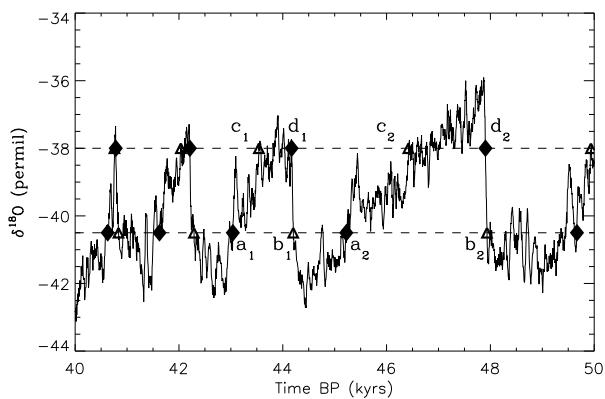


Figure 4:

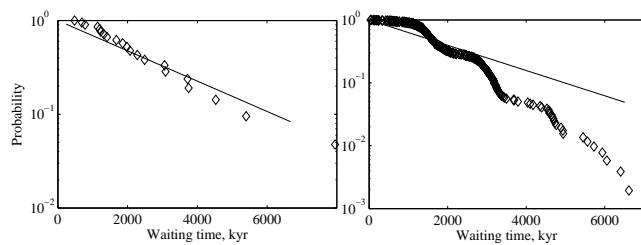


Figure 5:

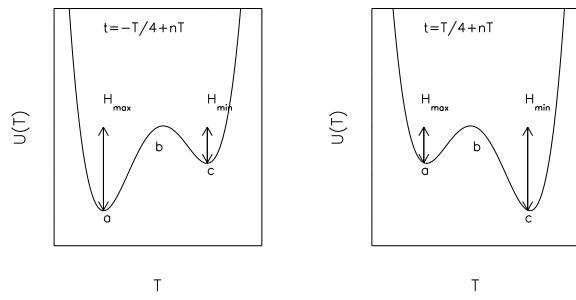


Figure 6:

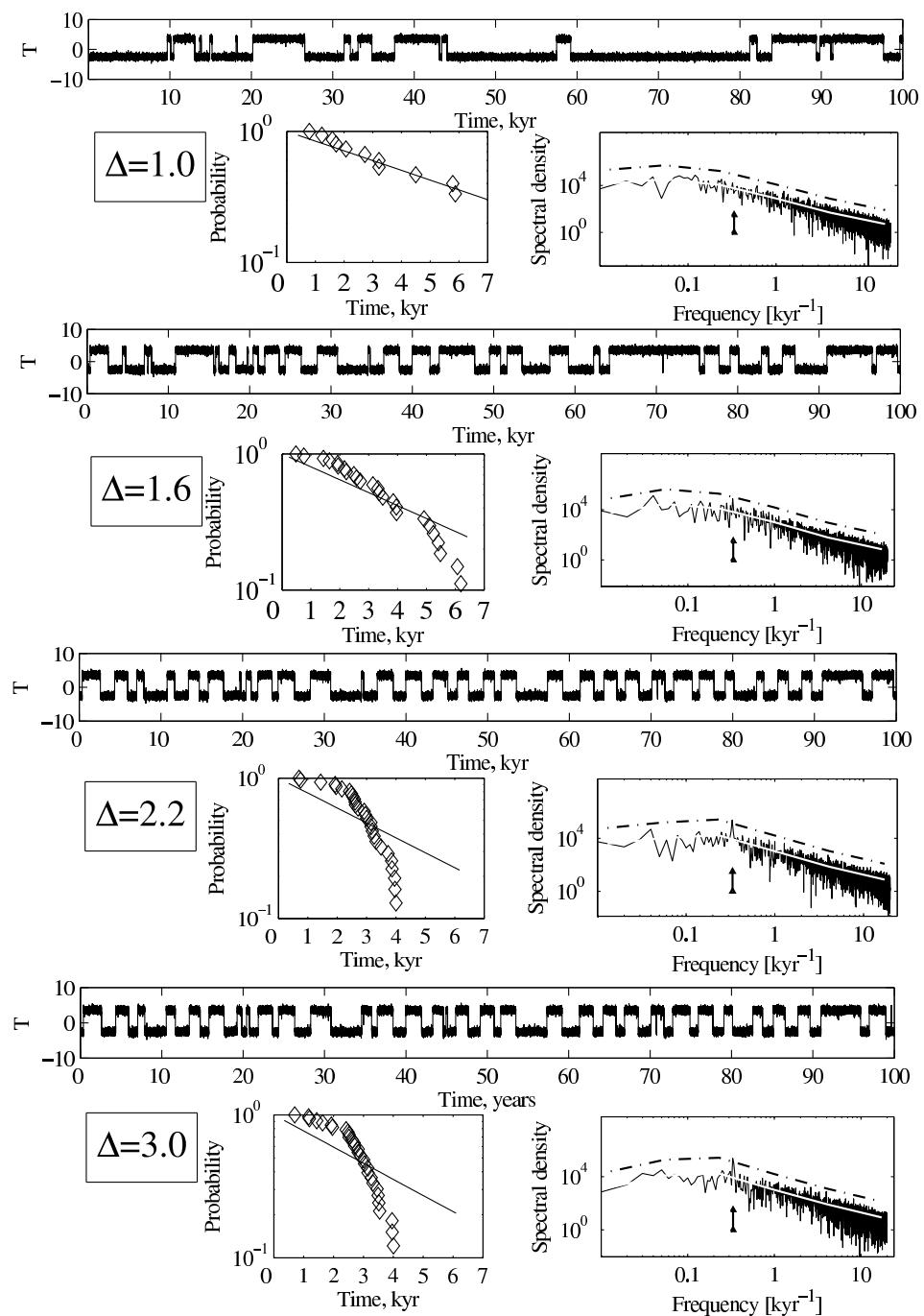


Figure 7:

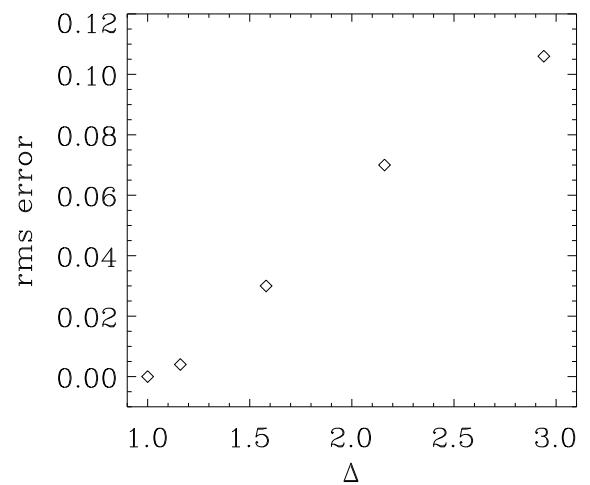


Figure 8:

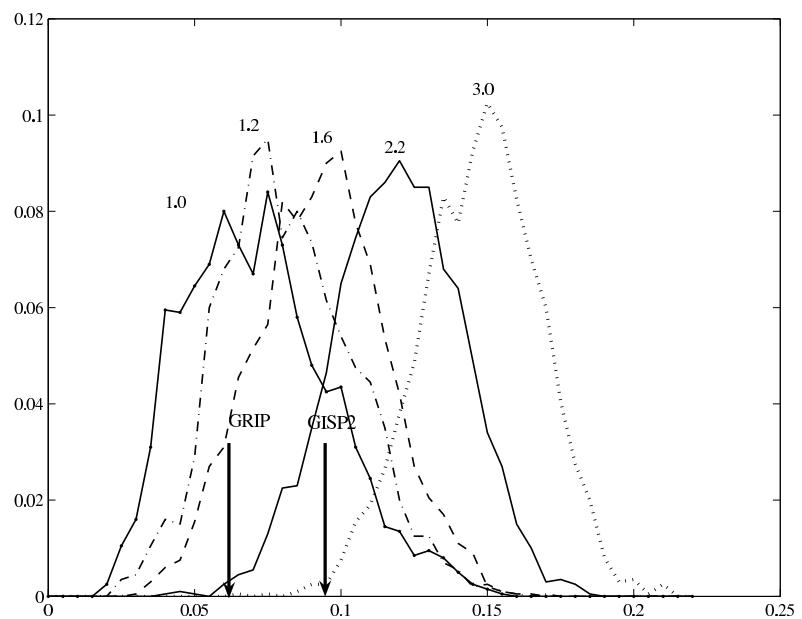


Figure 9: